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Nonlinear Bifurcation Analysis

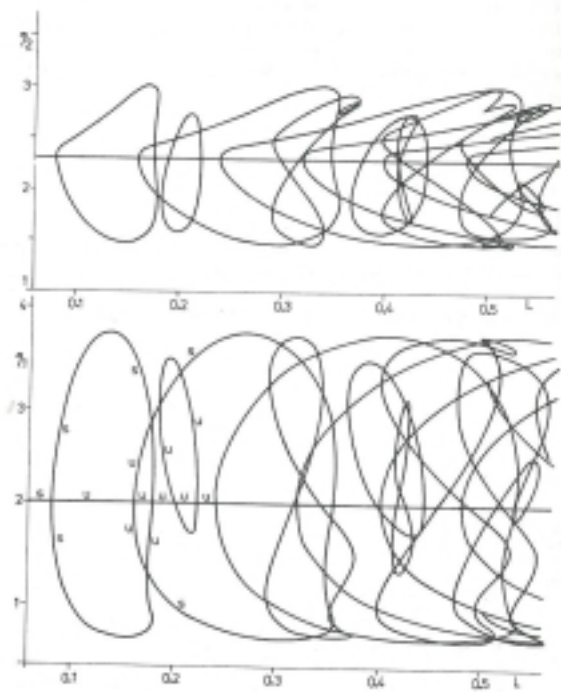


Figure 3.2 Dependence of the steady-state solutions on the parameter L ; Brusselator model, Example 6, Chapter 1, $A = 2$, $B = 4.6$, BC 2, $D_1 = 0.0016$, $D_2 = 0.008$, $v_1 = u(0)$, $v_2 = u(0)$, s —stable, u —unstable.

Bifurcation Analysis Levels

1. Locate

Find where B occurs while tracing a response

2. Determine subspace where "the action is"

Having located B , find y (particular solution) and z (homogeneous solution = null eigenvector in case of a simple bifurcation)


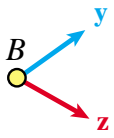
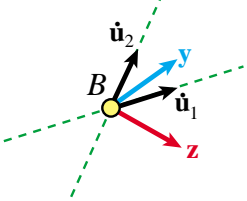
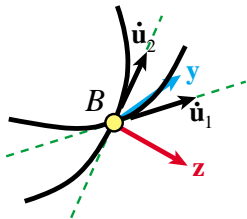
3. Branching analysis

Having located B and determined y and z , find the directions \dot{u}_1 and \dot{u}_2 of tangents to the equilibrium paths (branches) that pass through B . Requires analysis of $\ddot{r} = 0$

4. Branch curvature analysis

Having located B , and determined y , z , \dot{u}_1 and \dot{u}_2 find the curvatures of the equilibrium paths (branches) that pass through B . Requires analysis of $\ddot{r} = 0$

Bifurcation Analysis Levels: Pictures

Level	Diagram	Analysis Type	Rate Equation Order
1		Locate	1 (factorization)
2		Determine active subspace	1 (eigensolution)
3		Determine branch tangents at B	2 (if regular)
4		Determine branch curvatures at B	3 (if regular)

Can Commercial FEM Codes Do These Levels?

Level 1: yes

Level 2: yes

Level 3: very difficult

Level 4: impossible

Levels 3-4 can be directly done by hand, or with CAS help, only for **very simple problems.**

For more complicated ones, there are indirect ways to try such as injecting numeric perturbations.

Nonetheless, it **helps to know the theory to guide the computations (there is no such thing as a "black box" for this kind of analysis)**

Governing Equations Recap

Total residual equations:

$$\mathbf{r}(\mathbf{u}, \lambda) = \mathbf{0}$$

Parametric representation of the solutions (with t as pseudotime)

$$\mathbf{u} = \mathbf{u}(t) \quad \lambda = \lambda(t)$$

Two important special choices for t :

$$t = \lambda \quad t = s$$

which lead to **λ -parametrized** and **arclength** forms, respectively

Rate Forms

Rate forms of the governing residual equations are obtained by **repeated differentiation** of the total residual wrt pseudotime t :

$$\begin{aligned}\dot{\mathbf{r}} &= \mathbf{K}\dot{\mathbf{u}} - \mathbf{q}\dot{\lambda} = 0 \\ \ddot{\mathbf{r}} &= \mathbf{K}\ddot{\mathbf{u}} + \dot{\mathbf{K}}\dot{\mathbf{u}} - \dot{\mathbf{q}}\dot{\lambda} - \mathbf{q}\ddot{\lambda} = 0 \\ \dddot{\mathbf{r}} &= \mathbf{K}\dddot{\mathbf{u}} + \dot{\mathbf{K}}\ddot{\mathbf{u}} + \ddot{\mathbf{K}}\dot{\mathbf{u}} - \dot{\mathbf{q}}\ddot{\lambda} - \ddot{\mathbf{q}}\dot{\lambda} - \mathbf{q}\dddot{\lambda} = 0\end{aligned}$$

in which

$$\mathbf{K} = \frac{\partial \mathbf{r}}{\partial \mathbf{u}}, \quad \mathbf{q} = -\frac{\partial \mathbf{r}}{\partial \lambda} = -\mathbf{r}'$$

$\dot{\mathbf{r}} = 0$ is a system of **first-order rate equations**, also called the **incremental stiffness equations** or simply **stiffness equations**.
 $\ddot{\mathbf{r}} = 0$ is a system of **second-order rate equations**, also called the **stiffness rate equations**. And so on. For the branching analysis undertaken here we will go up to $\ddot{\mathbf{r}} = 0$.

Stiffness And Incremental Load Rates

In the second-order rate equation

$$\ddot{\mathbf{r}} = \mathbf{K}\ddot{\mathbf{u}} + \dot{\mathbf{K}}\dot{\mathbf{u}} - \dot{\mathbf{q}}\dot{\lambda} - \mathbf{q}\ddot{\lambda} = \mathbf{0}$$

the stiffness rate and incremental load arrays can be compactly expressed as

$$\dot{\mathbf{K}} = \mathbf{L}\dot{\mathbf{u}} + \mathbf{N}\dot{\lambda} \quad \dot{\mathbf{q}} = -(\mathbf{N}\dot{\mathbf{u}} + \mathbf{a}\dot{\lambda})$$

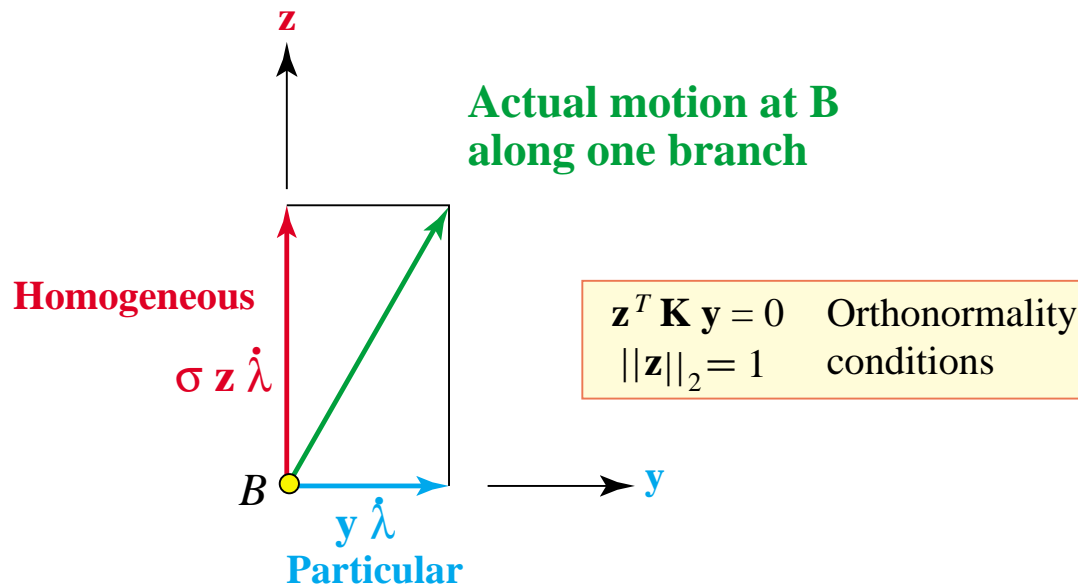
The entries of these matrices and vectors can be written in indicial form as

$$L_{ijk} = \frac{\partial^2 r_i}{\partial u_j \partial u_k} \quad N_{ij} = \frac{\partial^2 r_i}{\partial u_j \partial \lambda} = \frac{\partial K_{ij}}{\partial \lambda} \quad a_i = \frac{\partial^2 r_i}{\partial \lambda \partial \lambda} = \frac{\partial q_i}{\partial \lambda}$$

Note that \mathbf{L} is a **three-dimensional array** that may be called a **"cubic" matrix**. Postmultiplying a cubic matrix by a vector yields an ordinary matrix. For example $\mathbf{L} \dot{\mathbf{u}}$ in the expression of $\dot{\mathbf{K}}$ is a matrix

State Decomposition At Bifurcation

Recall from Chapter 28 the state decomposition at B



Finding the Null Eigenvector (Buckling Mode Shape)

To find \mathbf{z} , carry out an eigenvalue analysis of the tangent stiffness matrix \mathbf{K} at bifurcation point B . If this is a simple bifurcation point (which we will assume), \mathbf{z} is the **null eigenvector** of \mathbf{K} evaluated so that

$$\mathbf{K} \mathbf{z} = \mathbf{0}$$

This eigenvector is **guaranteed to have real entries since \mathbf{K} is symmetric**. **Normalize this vector to unit length** so that

$$\|\mathbf{z}\|_2 = 1$$

The **amplitude** of this eigenvector in the state rates near B is called σ .

Physically this eigenvector is often called the **buckling mode shape** but the name is not always accurate as regards physics

Control-State Rate Relation At Bifurcation

To carry out Level 2 analysis, recall from Chapter 28 that the **linearized relation** between control and state pseudotime rates measured from B can be expressed as a **linear combination** of \mathbf{y} and \mathbf{z} ,

$$\dot{\mathbf{u}} = (\mathbf{y} + \sigma \mathbf{z}) \dot{\lambda} \quad \mathbf{z}^T \mathbf{K} \mathbf{y} = \mathbf{y}^T \mathbf{K} \mathbf{z} = 0$$

Here \mathbf{z} is (as discussed in the previous slide) the unit-length-normalized null eigenvector at B (often called the buckling mode shape), whereas \mathbf{y} is the **particular solution** satisfying

$$\mathbf{K} \mathbf{y} = \mathbf{q}$$

as well as the orthonormality condition shown above; all quantities being evaluated at B

First-Order Rate Equations Do Not Cut It for Level 3

Can we do Level 3 branching analysis with $\dot{r} = 0$? This is the subject of HW Exercise 31.3. The analysis gives

$$\sigma = 0/0$$

Consequently the null eigenvector (buckling mode) amplitude is **indeterminate**, and the direction of the branches is **undefined**.

To get definite branching information it is necessary to proceed to the second-order rate equations. **Good news:** this will usually work (as long as the quadratic equation derived later does not fully degenerate). **Bad news:** we get more than one solution.

Level 3 Branching Analysis

To do Level 3 branching analysis we start from

$$\ddot{\mathbf{r}} = \mathbf{K}\ddot{\mathbf{u}} + \dot{\mathbf{K}}\dot{\mathbf{u}} - \dot{\mathbf{q}}\dot{\lambda} - \mathbf{q}\ddot{\lambda} = \mathbf{0}$$

Premultiply both sides by \mathbf{z}^T . Taking account of the bifurcation conditions

$$\mathbf{K}\mathbf{z} = \mathbf{z}^T\mathbf{K} = 0 \quad \mathbf{q}^T\mathbf{z} = \mathbf{z}^T\mathbf{q} = 0$$

we get

$$\mathbf{z}^T\dot{\mathbf{K}}\dot{\mathbf{u}} - \mathbf{z}^T\dot{\mathbf{q}}\dot{\lambda} = 0$$

Replacing the stiffness and incremental load rates by their compact expressions yields

$$\mathbf{z}^T(\mathbf{L}\dot{\mathbf{u}} + \mathbf{N}\dot{\lambda})\dot{\mathbf{u}} + \mathbf{z}^T(\mathbf{N}\dot{\mathbf{u}} + \mathbf{a}\dot{\lambda})\dot{\lambda} = 0$$

Level 3 Branching Analysis (cont'd)

Substituting the state rates by the state decomposition gives

$$\mathbf{z}^T [\mathbf{L}(\mathbf{y} + \sigma \mathbf{z})\dot{\lambda} + \mathbf{N}\dot{\lambda}] (\mathbf{y} + \sigma \mathbf{z})\dot{\lambda} + \mathbf{z}^T [\mathbf{N}(\mathbf{y} + \sigma \mathbf{z})\dot{\lambda} + \mathbf{a}\dot{\lambda}] \dot{\lambda} = 0$$

Collecting terms in the control rate and requesting a nontrivial solution gives the quadratic equation in σ

$$a\sigma^2 + 2b\sigma + c = 0$$

in which

$$a = \mathbf{z}^T \mathbf{L} \mathbf{z} \mathbf{z}$$

$$b = \mathbf{z}^T [\mathbf{L} \mathbf{z} \mathbf{y} + \mathbf{L} \mathbf{y} \mathbf{z} + 2\mathbf{N} \mathbf{z}]$$

$$c = \mathbf{z}^T [\mathbf{L} \mathbf{y} \mathbf{y} + 2\mathbf{N} \mathbf{y} + \mathbf{a}]$$

Level 3 Branching Analysis (cont'd)

This quadratic equation generally provides two roots: σ_1 and σ_2 . In what follows **we shall assume that the roots are real** (see Chapter for discussion) . Substitution into

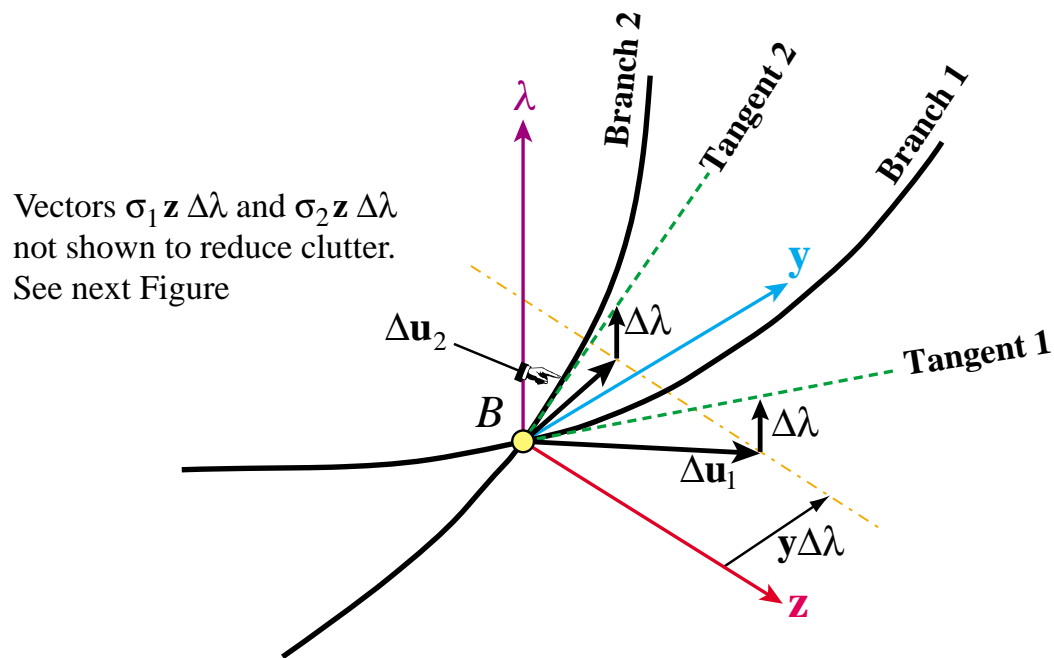
$$\dot{\mathbf{u}} = (\mathbf{y} + \sigma \mathbf{z}) \dot{\lambda}$$

furnishes the two **branch directions** at the bifurcation point:

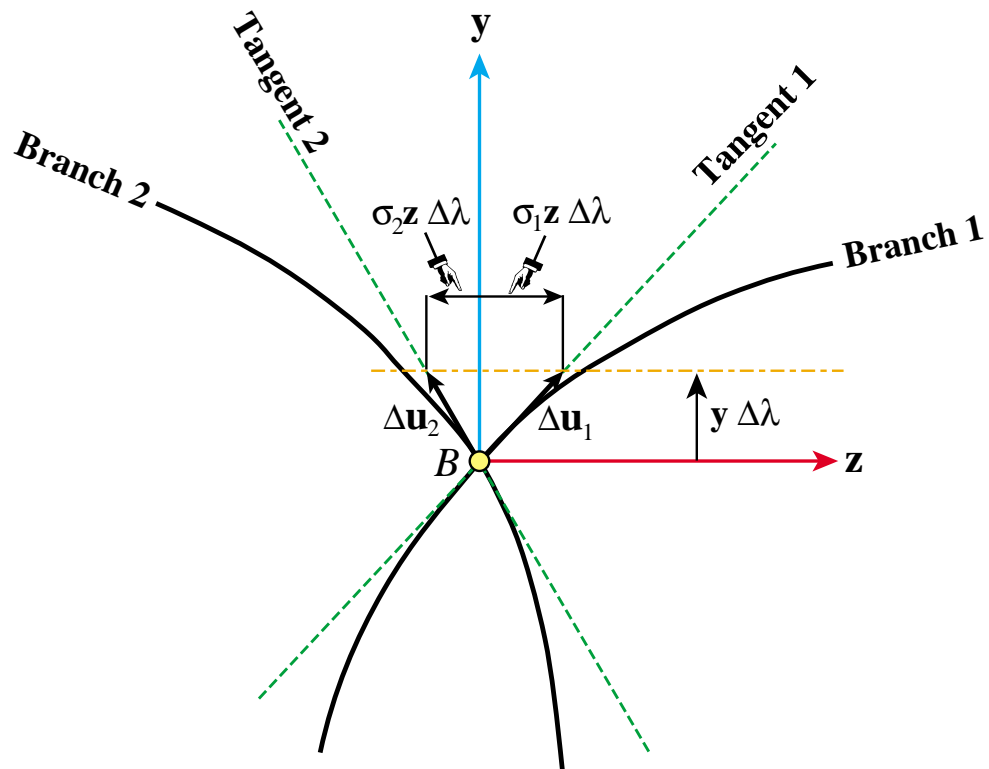
$$\dot{\mathbf{u}} = (\mathbf{y} + \sigma_1 \mathbf{z}) \dot{\lambda} \quad \dot{\mathbf{u}} = (\mathbf{y} + \sigma_2 \mathbf{z}) \dot{\lambda}$$

These are sketched in the next two slides, first in the three dimensional space $(\mathbf{y}, \mathbf{z}, \lambda)$, then projected on the (\mathbf{y}, \mathbf{z}) plane.

Branching Directions In (y, z, λ) Space



Branching Directions On (y, z) Plane



Special Forms Of The σ -Quadratic Equation

(questions to class - See Chapter for discussion)

$$a\sigma^2 + 2b\sigma + c = 0$$

What happens if $a = 0$?

What happens if $a = c = 0$?

Is $a = b = 0$ possible?

What happens if $a = b = c = 0$?

What Happens At A Multiple (a.k.a. Compound) Bifurcation Point?

Suppose now that **two bifurcation points coincide**. There are now **two independent null eigenvectors**:

$$\mathbf{z}_1 \quad \mathbf{z}_2$$

which together **span a null eigenspace** Z of dimension two (in linear algebra terminology, the kernel space of K evaluated at bifurcation point B has dimension 2)

Assuming all roots of the branching equations are real, **how many equilibrium branches** go through B ?

Unexpected Answer

If the bifurcation point has multiplicity M , the number of branches crossing there ranges from

$$2M-1 \quad \text{to} \quad (3^M - 1)/2$$

depending on local symmetries. Results were published by Sewell (1968), Johns and Chilver (1971) for general Jacobians (not necessarily derivable from potentials)

The complexity of the analysis **grows enormously** as M increases:

For $M=2$ the range is **3 to 4**

For $M=5$ (NASA benchmark: composite panel with a rectangular cutout) the range is **31 to 121**